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Wave Journal Bearings Under Dynamic Loads

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WAVE JOURNAL BEARINGS UNDER DYNAMIC LOADS

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SUMMARY

The dynamic behavior of the wave journal bearing was determined by running a three-wave bearing with an eccentrically mounted shaft. A transient analysis was developed and used to predict numerical data for the experimental cases. The three-wave journal bearing ran stably under dynamic loads with orbits well inside the bearing clearance. The orbits were almost circular and nearly free of the influence of, but dynamically dependent on, bearing wave shape.

Experimental observations for both the absolute bearing-housing-center orbits and the relative bearing-housing-center-to-shaft-center orbits agreed well with the predictions. Moreover, the subsynchronous whirl motion generated by the fluid film was found experimentally and predicted theoretically for certain speeds.

SYMBOLS

B	damping coefficient between bearing housing and its support, N·s/m
$B_{ij(i = r,t;j = r,t)}$	film damping coefficients, N·s/m
C	bearing radial clearance, m
F_r, F_t	fluid film force components along and perpendicular to line of centers $O_1 - O$ (fig. 1), N
$F_{i0(i = r,t)}$	steady-state fluid film force components along and perpendicular to line of centers $O_1 - O$, N
F_ζ, F_η	fluid film force components along and perpendicular to line $O_0 - O$ (fig. 1), N
h	fluid film thickness, m
K	stiffness coefficient between bearing housing and its support, N/m
$K_{ij(i = r,t;j = r,t)}$	film stiffness coefficients, N/m
M	bearing-housing total mass, kg
O	bearing-housing center (fig. 1)
O_0	fixed center of rotation (fig. 1)
O_1	shaft center (fig. 1)
p	fluid film pressure, Pa
R	shaft radius (bearing normal radius), m
r	coordinate along line of centers
S	space (displacement) between shaft center O_1 and bearing-housing center O, m
$S_{i(i = r,t)}$	S components along and perpendicular to line of centers $O_1 - O$, m
t	time, s; coordinate perpendicular to line of centers
V	velocity between shaft center O_1 and bearing-housing center O, m/s

$V_{i(i = r,t)}$	V components along and perpendicular to line of centers $O_1 - O$, m/s
V_n	difference between shaft surface and bearing surface speed projected on perpendicular direction to shaft surface, m/s
V_θ	component of shaft surface speed along its circumference, m/s
z	axial coordinate parallel to shaft axis
ε	bearing eccentricity relative to shaft (relative movement) (fig. 1), $O_1 - O$, m, also bearing line of centers
ε/C	bearing sleeve-to-shaft eccentricity ratio
ε_0	bearing run-out (absolute movement) (fig. 1), $O_0 - O$, m
ε_0/C	bearing-housing absolute eccentricity ratio, $\varepsilon_0 = O_0 - O$
ζ, η	axes along and perpendicular to direction $O_0 - O$ (fig. 1)
θ	angular coordinate along shaft circumference, rad
μ	fluid film dynamic viscosity, $N \cdot s/m^2$
ρ	fixed shaft run-out (fig. 1), $O_0 - O_1$, m
φ	angle between line of centers $O_1 - O$ and $O_0 - O$ (fig. 1), rad
ψ	rotation angle of $O_0 - O$ around O_0 (fig. 1), rad
Ω	rotation angle of $O_0 - O_1$ around O_0 (fig. 1), ωt, rad
ω	angular rotation speed (fig. 1), rad/s

INTRODUCTION

The wave bearing concept has been under development since 1992. Thus, the steady-state and dynamic performance under fixed side load (refs. 1 and 2) and the influence of both the number of waves and the ratio of wave amplitude to radial clearance (refs. 3 and 4) have been analyzed. Moreover, the steady-state characteristics of the wave journal bearing and its dynamic stability have been experimentally measured. Good agreement was found between the experimental data and theoretical predictions (refs. 5 to 8). In addition, the experimental work revealed good dynamic behavior of the wave bearing when subsynchronous whirl motion occurred. The wave bearing performed well, keeping the orbit of the subsynchronous motion inside the bearing clearance (refs. 7 and 8). Consequently, the wave bearing should perform well under the dynamic loading conditions that often occur in most rotating machinery. Any rotor can be subject to a dynamic load caused by an unbalance, or a run-out, of the shaft. This dynamic load is a rotating load that has a rotational speed equal to the rotor speed. Such a load can be simulated by running the bearing with a shaft that has a fixed run-out. Therefore, a transient analysis was performed to predict bearing behavior under a rotating load. Then an experiment was conducted to record the orbits of the bearing-housing center when the shaft has a known fixed run-out.

ANALYSIS

Bearing-housing-center movement can be studied by using the motion equation of the center along and perpendicular to the radial direction $O_0 - O$ (axes ζ and η in fig. 1):

$$MC\left[\frac{d^2\varepsilon_0}{dt^2}-\varepsilon_0\left(\frac{d\psi}{dt}\right)^2\right]+K\varepsilon_0+B\frac{d\varepsilon_0}{dt}=F_\zeta$$

$$MC\left(\varepsilon_0\frac{d^2\psi}{dt^2}+2\frac{d\psi}{dt}\frac{d\varepsilon_0}{dt}\right)+K\varepsilon_0\psi+B\varepsilon_0\frac{d\psi}{dt}=F_\eta \tag{1}$$

Also, figure 1 shows that the eccentricity $\varepsilon = O_1 - O$ (where $O_1 - O$ is the line joining the shaft center O_1 and the bearing-housing center O) and that the shaft run-out $\rho = O_0 - O_1$. Assuming that the motion starts from the downward vertical where the shaft and the bearing are concentric ($\varepsilon = 0$), when the shaft rotates around O_0 with the angular speed ω, ρ makes the angle Ω and drives the bearing so that ε_0 makes the angle ψ.

The governing equations (1) are two scalar, coupled, nonlinear ordinary differential equations. These equations are integrated simultaneously by using a fourth-order Runge-Kutta method for known values of M, C, F_ζ, F_η, K, and B and initial values of ε_0, ψ, $d\varepsilon_0/dt$, and $d\psi/dt$ (ref. 9). The fluid film forces applied to the bearing surface are

$$F_\zeta = F_r\cos\varphi + F_t\sin\varphi$$

$$F_\eta = F_r\sin\varphi - F_t\cos\varphi \tag{2}$$

The projections of fluid film force along and perpendicular to the line of centers are

$$F_r = F_{r_0} + K_{rr}S_r + K_{rt}S_t + B_{rr}V_r + B_{rt}V_t$$

$$F_t = F_{t_0} + K_{tr}S_r + K_{tt}S_t + B_{tr}V_r + B_{tt}V_t \tag{3}$$

The bearing steady-state force and dynamic stiffness and damping coefficients can be computed by integrating the Reynolds pressure equation at each time step location of the shaft with respect to the bearing. This equation, assuming the gas will expand isothermally, is

$$\frac{\partial}{R\partial\theta}\left(\frac{h^3}{\mu}p\frac{\partial p}{R\partial\theta}\right)+\frac{\partial}{\partial z}\left(\frac{h^3}{\mu}p\frac{\partial p}{\partial z}\right)=6\left[2pV_n+2\frac{\partial(ph)}{\partial t}+pV_\theta\frac{\partial h}{R\partial\theta}+h\frac{\partial(pV_\theta)}{R\partial\theta}\right] \tag{4}$$

The Reynolds equation (4) can be integrated by using its complex form and a small perturbation technique. This procedure is described, for instance, in reference 10.

The solution procedure can start with an input data set (bearing length, diameter, radial clearance, shaft turning speed, shaft run-out, and the time step). In addition, a set of starting values at time = 0 are required:

$$\varepsilon_0 = \rho, \quad \frac{\partial\varepsilon_0}{\partial t} = 0$$

$$\psi = 0, \quad \frac{\partial\psi}{\partial t} = 0 \tag{5}$$

Then, at each time step, where ε_0, ψ, and Ω ($\Omega = \omega t$) are known, the O_0O_1O triangle (fig. 1) is known, and all geometrical parameters as well as displacements and velocities can be calculated. Therefore, the Reynolds equation (4) can be integrated over the fluid film. Then, all parameters of the motion equation (1) are known as well as the starting values for the next time step (ε_0, ψ, and their time derivatives Ω). The procedure is repeated until the orbits are completed.

APPARATUS

The wave bearing rig described in references 5 to 8 was used to perform the experimental work. The axis of the spindle that drives this rig is vertical, and the experimental bearing housing is mounted on the rig table and supported by two pressurized thrust plates. This configuration keeps the bearing housing stiff in the axial and angular directions but allows it to move freely in the radial direction. The experimental shaft is an extension of the rig spindle shaft. It is mounted into the tapered end of the spindle shaft with a fixed run-out (for this experiment 11±0.1 μm). A cross section by a horizontal plane of the experimental bearing is shown in figure 2. The fixed rotation center for the system is O_0. The centers of the shaft and the bearing housing are O_1 and O, respectively. The shaft run-out $O_0 - O_1$ is fixed.

The goal of this work was to record and predict the absolute and relative orbits of the bearing-housing center O. The motion of the center O can be observed like an absolute motion for instance with regard to the center of rotation O_0 or like a relative motion with regard to the center of the shaft O_1. Figure 3 shows the experimental bearing setup. Two sets of light-beam proximity probes were used. Two probes were located at 90° in the bottom side of the bearing housing and "looking" at the shaft. These probes detected the orbit of the bearing-housing center relative to the shaft center ($O - O_1$). The second set of two probes were located also at 90° but held by supports fixed on the rig table and "looking" at the bearing housing. These latter probes detected the absolute orbit of the bearing-housing center ($O - O_0$). A polished circumferential strip was made on the outside bearing-housing surface to avoid asperity noise from its roughness. The light-beam probes were calibrated by using the known fixed run-out of the shaft. The displacement of the shaft was measured with a precision of 0.1 μm. The theoretical predictions of the orbits were made through a transient analysis of the bearing-housing-center motion.

RESULTS AND DISCUSSION

The experimental bearing was 51±0.01 mm in diameter, 58±0.01 mm in length, 20±1 μm in radial clearance, and 2.2±0.01 kg in mass. The bearing had three waves with a 0.5±0.07 ratio of wave amplitude to radial clearance. The shaft was set with an 11±0.1 μm run-out. The damping, B in eqs. (1), in the bearing-housing support and connection system was found to be 0.05 N·s/m. The stiffness, K in eqs. (1), had little influence on the bearing orbits and was approximately zero. The top proximity probes (fig. 3) produced 500 mV for 5.78- and 4.78-μm displacements in the horizontal and vertical directions, respectively, and the bottom probes produced 500 mV for 6.11- and 6.90-μm displacements in the horizontal and vertical directions, respectively. (Horizontal and vertical directions refer to the directions on the oscilloscope photographs shown on the right sides of figures 4 and 5, 90° apart in the physical plane.

The test rig was run at four speeds up to 5540 rpm. Below 3100 rpm both the observed and predicted orbits of the bearing-housing center showed that a subsynchronous whirl motion took place inside the bearing clearance. Figure 4 shows both the predicted and observed orbits for relative and absolute motion of the bearing-housing center when the shaft rotated at 2156 rpm. When the speed increased above 3100 rpm, the motion stabilized, as shown in figure 5 for a shaft speed of 5539 rpm.

Both the absolute and relative observed orbits of the bearing-housing center are shown as oscilloscope photographs on the right sides of figures 4 and 5. On the left sides of these figures the computed orbits are presented with a time step of 0.000001 s (10 μs) and for 30 000 steps. The experimental orbits appeared as ellipses rather than circles because of the difference in the probe sensitivity in the horizontal and vertical directions mentioned above. Both experimental orbits in figure 4 have a specific pattern caused by subsynchronous whirl motion. The transient analysis also revealed this pattern. Both the experimental and theoretical absolute orbits (fig. 4(a)) were within a radius of 5 to 12 μm. Both the experimental and theoretical relative orbits (fig. 4(b)) were within a radius of approximately 5 μm.

The bearing stability increased as the running speed of the rig increased. Figure 5 shows the results for 5539 rpm. The experimental orbits were perfectly stable. The shaft run-out made large absolute orbits of the bearing housing (right side of fig. 5(a)). However, the radius of the relative orbits was approximately 2.5 μm (right side of fig. 5(b)) despite the 11±0.1 μm shaft run-out (i.e., the bearing followed the shaft very well). The predicted orbits, shown on the left side of figure 5, matched very well with the observed orbits. The theory also showed that the bearing would run stably. After a couple of rotations from the starting point the orbits were stable, keeping almost the same path.

The relative orbits of the bearing housing increased but the absolute orbits decreased as the speed increased. This effect showed the influence of both external damping and bearing inertia on the magnitude of orbit radius. In addition, the bearing actually ran more and more stably as speed increased, and the theory showed that the number of rotations before the bearing reached a stable orbit would diminish as speed increased.

All runs showed only a small influence of the bearing wave shape on the orbit shape despite the experimental bearing's large wave amplitude ratio, 0.5±0.07. This result confirmed that a wave bearing with few waves, such as three, worked well under dynamic loads. The bearing behaved in such a way as to average the influence of the waves.

Two types of shaft-centered motion can be defined with respect to the center of the bearing: (i) stable unbalance or run-out orbits (e.g., fig. 5(b)), where the center of unbalance rotates at shaft frequency; and (ii) fluid-film-induced unstable whirl orbits (e.g., fig. 4(b)) that are superimposed over the stable unbalance or run-out orbits at a specific frequency different from the rotation frequency. The unbalance motion (i) is seen in each graph, but the unstable whirl (ii) occurs only at specific rotational speeds.

CONCLUSIONS

The dynamic behavior of the wave journal bearing was determined by running a three-wave bearing with an eccentrically mounted shaft. The following conclusions were reached:

1. A dynamically loaded three-wave journal bearing can run stably, averaging its behavior when the wave exposure to the load is changing. The orbit radius of the relative motion between the shaft and the sleeve is smaller than the bearing clearance, and the motion is contained within the bearing clearance. The orbits are almost circular and nearly free of, but dynamically dependent on, the influence of bearing wave shape.
2. Good agreement between experimentally observed and theoretically predicted orbits was found at all tested speeds for both relative and absolute motions.
3. The subsynchronous whirl motion influences the bearing-housing-center orbits if the bearing speeds are in the region where the bearing itself is susceptible to subsynchronous whirl instability. When the bearing runs under such circumstances, the orbits show a specific pattern. This pattern was observed experimentally and was also confirmed theoretically by the transient analysis.

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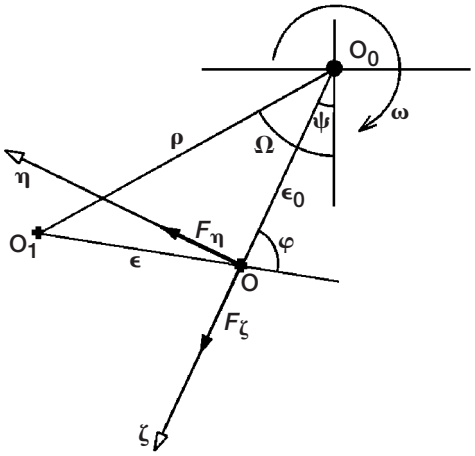

Figure 1.—Geometry of three-wave journal bearing. Rotation center, O_0; shaft center, O_1; and bearing-housing center, O.

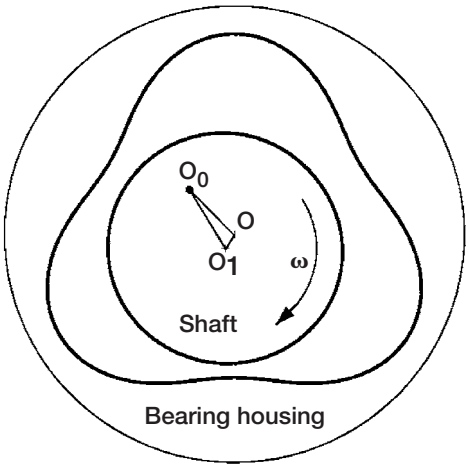

Figure 2.—Cross section of three-wave journal bearing by horizontal plane.

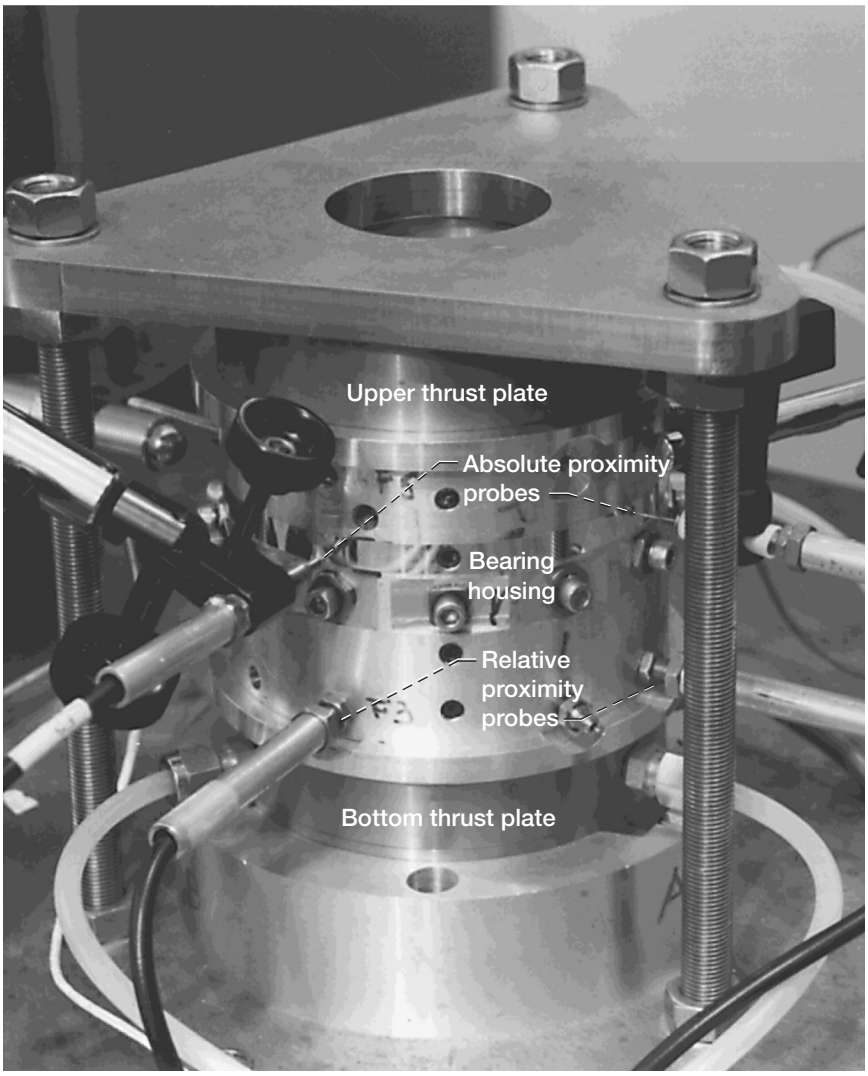

Figure 3.—Experimental bearing setup.

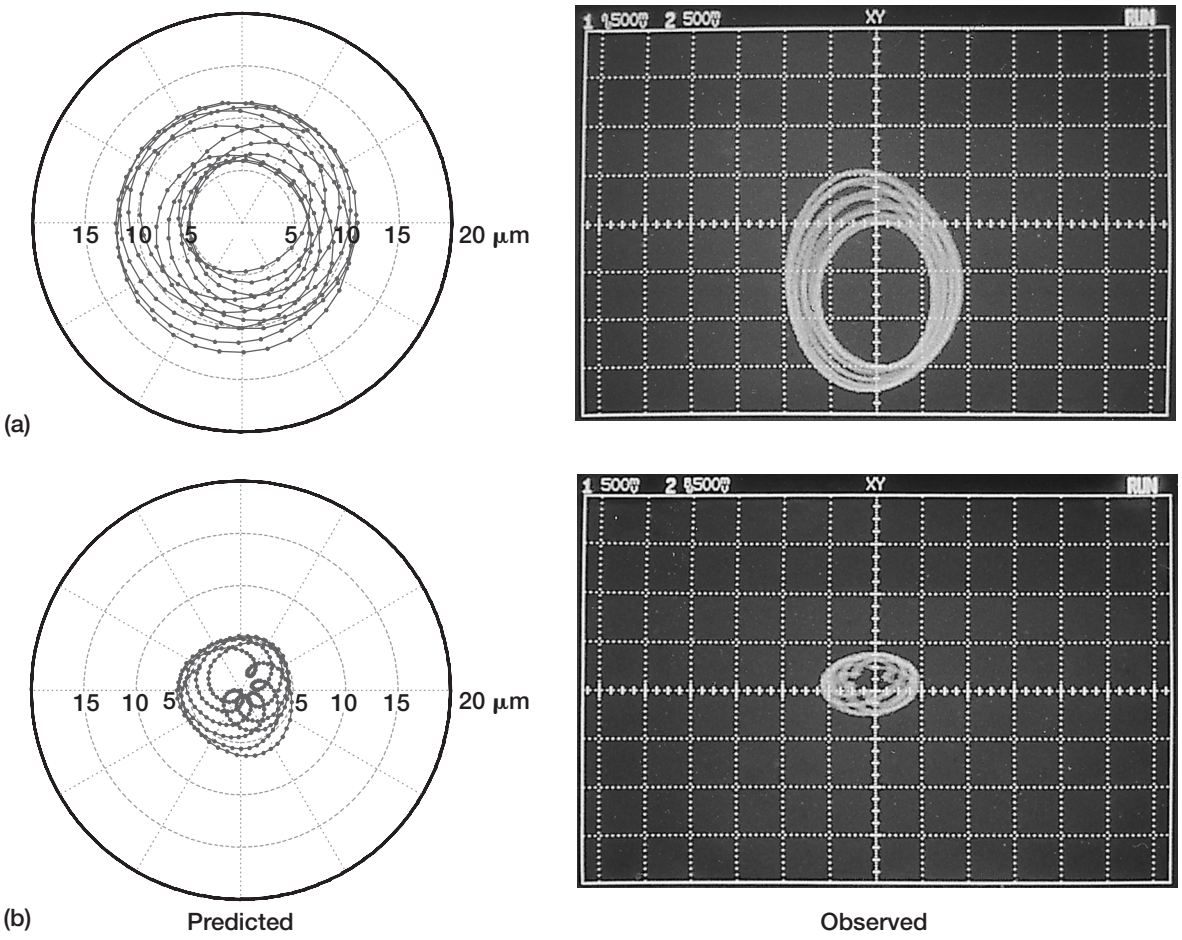

Figure 4.—Predicted and experimentally observed orbits of three-wave journal bearing at 2156-rpm shaft rotating speed. (a) Absolute bearing-housing-center orbits. (b) Relative bearing-housing-center-to-shaft-center orbits.

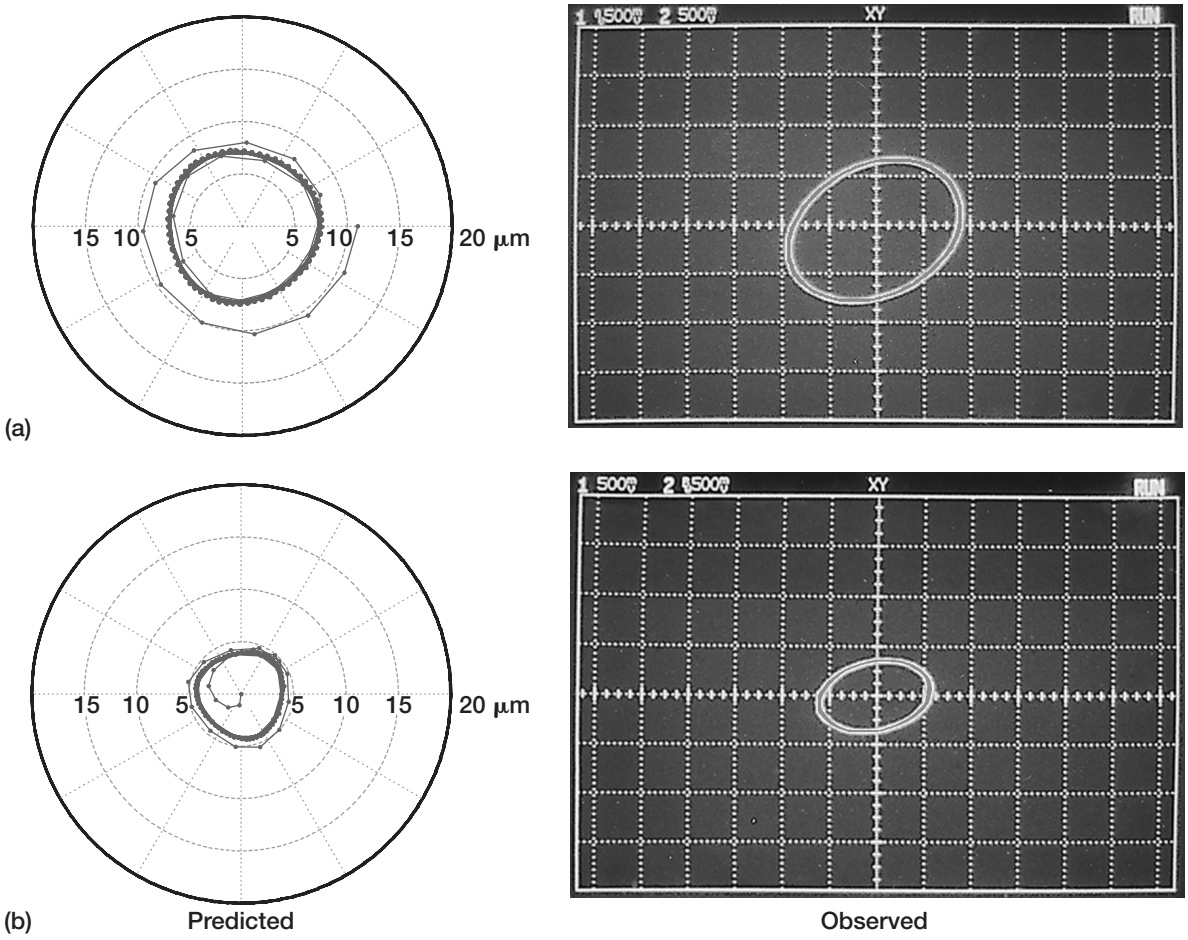

Figure 5.—Predicted and experimentally observed orbits of three-wave journal bearing at 5539-rpm shaft rotating speed. (a) Absolute bearing-housing-center orbits. (b) Relative bearing-housing-center-to-shaft-center orbits.

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13. ABSTRACT ***(Maximum 200 words)***

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